

# Detection Of Ionizing Radiation



# Agenda

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- Detection of ionizing radiation (photons and charged particles)
  - Solid-state detectors (Ge, Si)
  - Gas amplification detectors (Ionization chamber, proportional counter, Geiger counter)

Reading: Knoll Ch.12.I-12.IV

- Phenomenological model of matter ionization by particles
  - Electronic stopping
  - Bethe-Bloch Formula
  - Examples
  - Range and specific ionization
  - Stopping power curves, energy loss in thin foils

Reading: Knoll Ch. 6.I-6.V

- Spurious response of gas counters to photons

# Detector Design Principles

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## Ionization (charge separation) Detectors

- Ionization chambers (solid-state and gas)
- Gas Amplification Dets
  - Proportional counters
  - Avalanche counters
  - Geiger-Müller counters
- Cloud/bubble chambers
- Solid track detectors

## Scintillation Detectors

- Phosphorescence counters
- Fluorescence counters (inorganic solid crystal scintillators, organic solid and liquid scintillators)
- Čerenkov counters

## Associated Techniques

- Photo sensors and multipliers
- Charged-coupled devices
- Electronic pulse shape analysis
- Processing/acquisition electronics

# Ionization Chambers (Solid-State and Gas Medium)

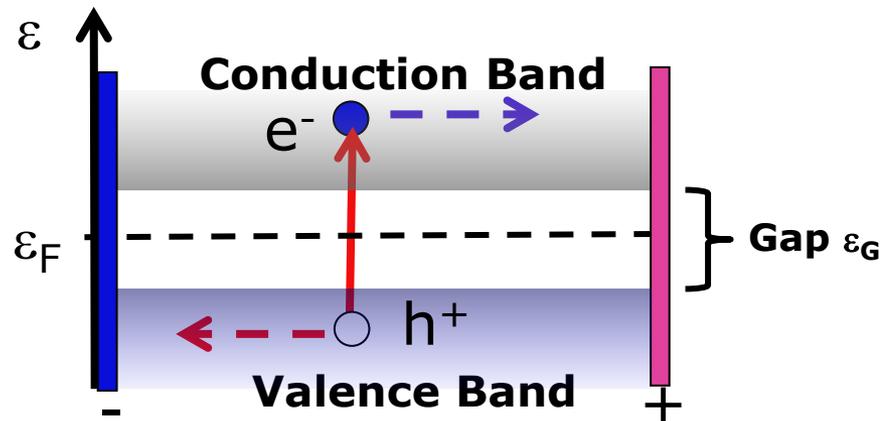
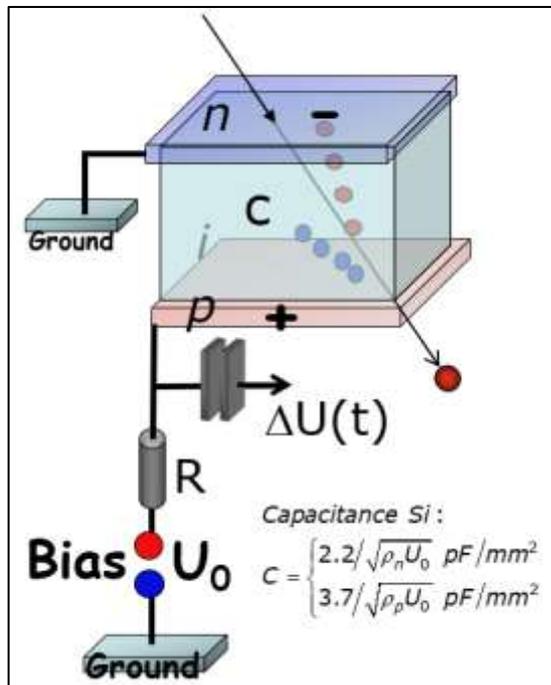
**General principle:** Radiation dissipates energy  $E$  via production of electron-ion ( $e^-$ ,  $h^+$ ) pairs in a medium enclosed between electrodes (Anode, Cathode). Electronic  $E$  signal picked up at A or C.

Gas volume between capacitor C electrodes.

Energy  $E \rightarrow N_{\text{ion pairs}} = E/\varepsilon_{\text{ip}}(\text{gas})$

Semiconductor  $n$ -,  $p$ -,  $i$ -types  $Si$ ,  $Ge$ ,  $GaAs$ ,...

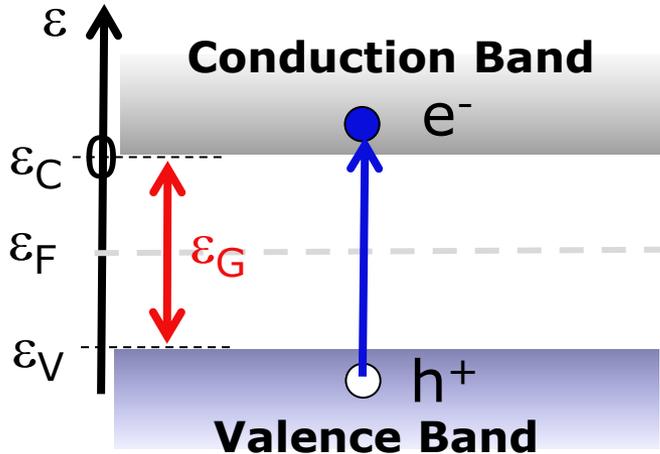
Band structure of solids VB gap CB.



Ionization lifts  $e^-$  up to CB, leaves hole  $h^+$  in VB  $\rightarrow$  free charge carriers, produce  $\Delta U(t) \sim E$ .

# Particles and Holes in Hyper **Pure** Semi-Conductors

Fermi gas of electrons (and holes)  
Fermion statistics @ temperature  $T$ :



$$e^- : f_e(\varepsilon) = \left[ 1 + \exp\left(\frac{\varepsilon + \varepsilon_G/2}{kT}\right) \right]^{-1}$$

$$h^+ : f_h(\varepsilon) = \left[ 1 + \exp\left(\frac{-\varepsilon + \varepsilon_G/2}{kT}\right) \right]^{-1}$$

Small gaps  $\varepsilon_G$  (Ge)  $\rightarrow$   
high thermal currents.  
Reduce by cooling.

$n_e, n_h = \#$  of occupied  $e^-$  or  $h^+$  states  
 $f_e, f_h \leq 1$  occupation numbers

$$n_e(\varepsilon) = \frac{(2m)^{2/3} V}{2\pi^2 \hbar^3} \sqrt{\varepsilon} \cdot f_e(\varepsilon) \quad V = \text{volume}$$

$$n_h(\varepsilon) = \frac{(2m)^{2/3} V}{2\pi^2 \hbar^3} \sqrt{|\varepsilon|} \cdot f_h(\varepsilon) \quad n_e = n_h !!$$

$$\varepsilon_F = \varepsilon_C - \varepsilon_G/2 = -\varepsilon_G/2 \quad \text{for } \varepsilon_C := 0$$

$$f_e(\varepsilon) = \left[ 1 + \exp\left(\frac{\varepsilon - \varepsilon_F}{kT}\right) \right]^{-1}$$

$$\xrightarrow{kT \approx 25 \text{ meV} \ll \varepsilon_G} \exp\left(-\frac{\varepsilon + \varepsilon_G/2}{kT}\right)$$

$$\langle n_e^2 \rangle = \langle n_e n_h \rangle = \left( \frac{(2m)^{2/3} V}{2\pi^2 \hbar^3} \right)^2 \langle \varepsilon \rangle \exp\left(-\frac{\varepsilon_G}{kT}\right)$$

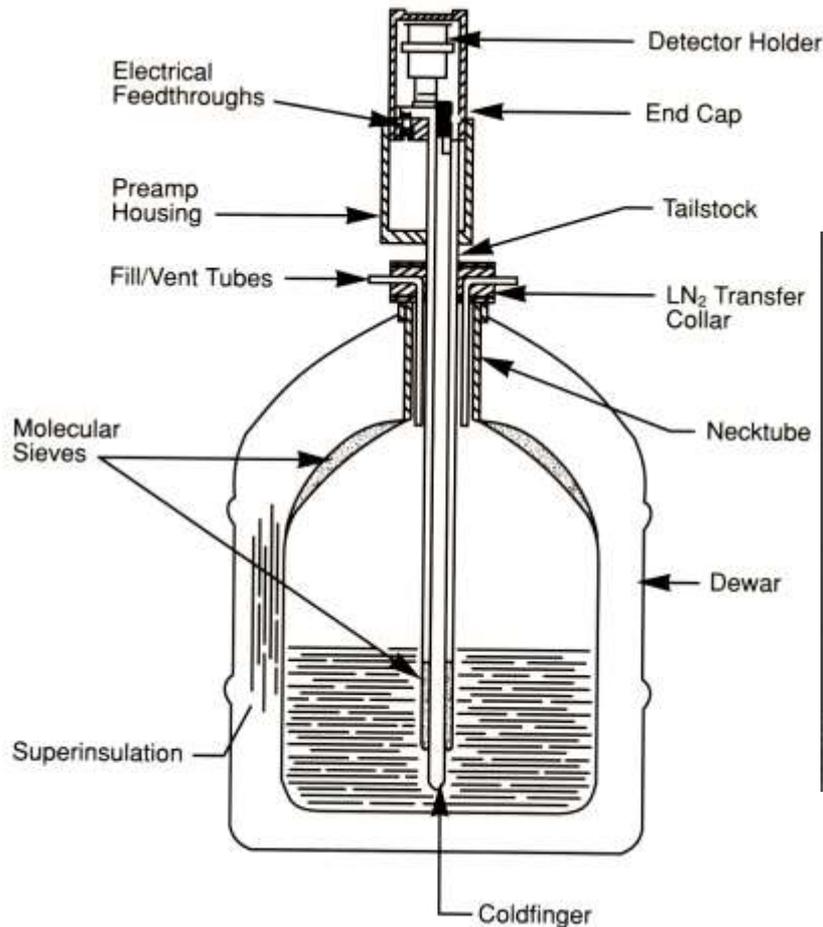
$$\langle n_e \rangle_{rms} \sim \exp\left(-\frac{\varepsilon_G}{2kT}\right)$$

$\propto$  noise generating  
conductivity at  $T$

# Hyperpure (Intrinsic) Ge $\gamma$ -ray Detectors

Hyper-pure Ge detectors for  $\gamma$ -rays use because of small gap  $E_G$ , cool to  $-77^\circ\text{C}$  ( $\text{LN}_2$ ). Simple band structure.

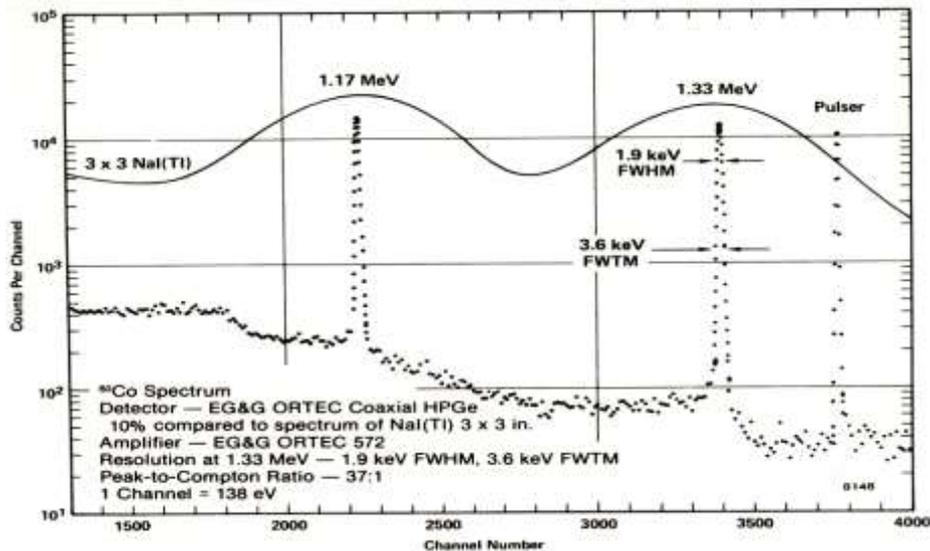
## Ge Cryostat (Canberra)



## Ge cryostat geometries (Canberra)



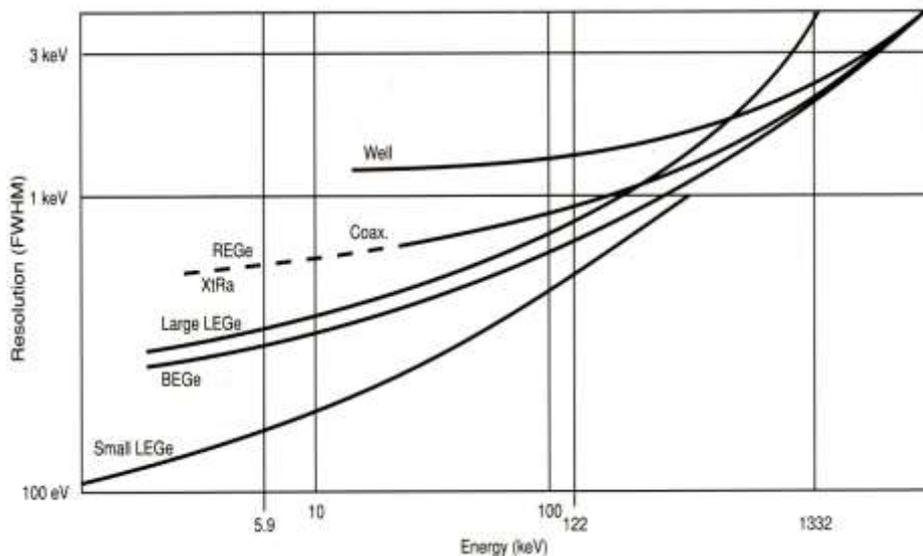
# Properties of Ge Detectors: Energy Resolution



Superior energy resolution, compared to NaI

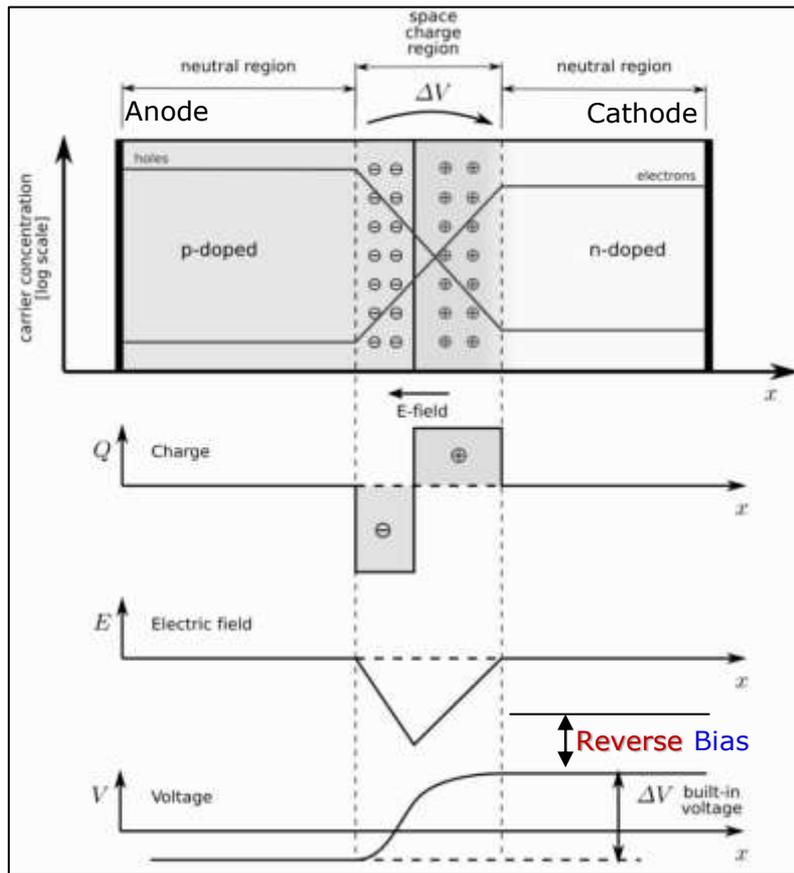
$$\Delta E_{\gamma} \sim 0.5 \text{ keV} @ E_{\gamma} = 100 \text{ keV}$$

Higher peak/Compton ratios



Size=dependent small detection efficiencies of Ge detectors  $\epsilon \sim 10\% \rightarrow$  solution: bundle in  $4\pi$ -arrays *GammaSphere, Greta EuroGam, Tessa, ...*

# Alternative: Semiconductor Junctions and Barriers



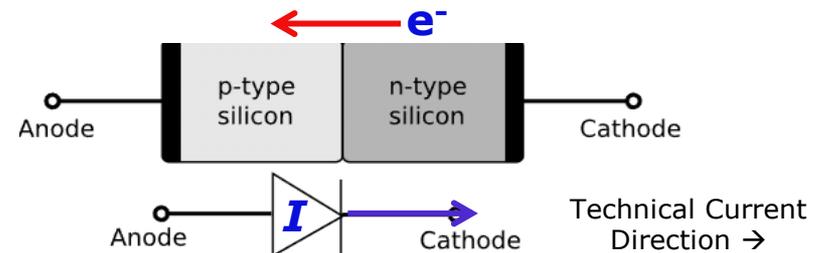
Need detector for rad-induced charges  
 → otherwise, no free carriers allowed.

Difficult to make: perfect *i*-type (intrinsic) Si  
 = chemical Group IV.

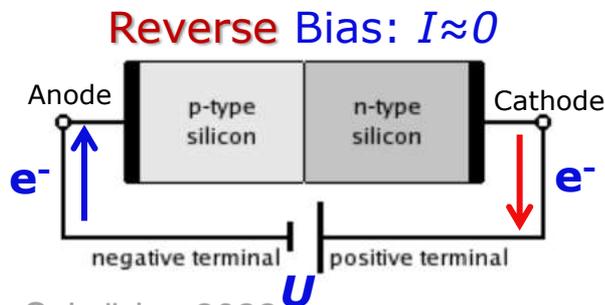
**Trick:** to make fully depleted Si → SC junction  
*n*-type Si: by doping with *Li* or Group V  
 e<sup>-</sup> donor atoms (*P, Sb, As*),  
*p*-type Si: by doping with Group III  
 e<sup>-</sup> acceptor atoms (*B, Al, ..*).

Junctions diffuse donors and acceptors into Si  
 bloc from different ends → interface → e<sup>-</sup>/h<sup>+</sup>  
 annihilation → space charge = depleted zone

## Semiconductor Diode

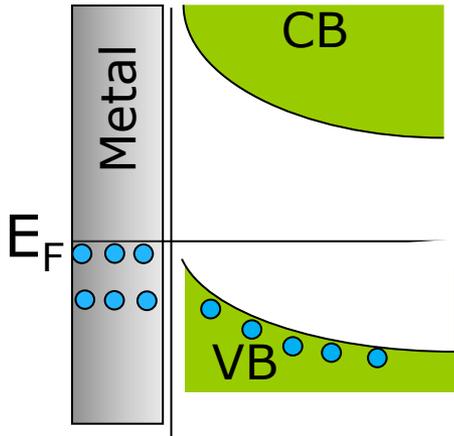


Electrons move easily through the junction *from n to p*  
 but *not from p to n*, and the reverse is true for holes.



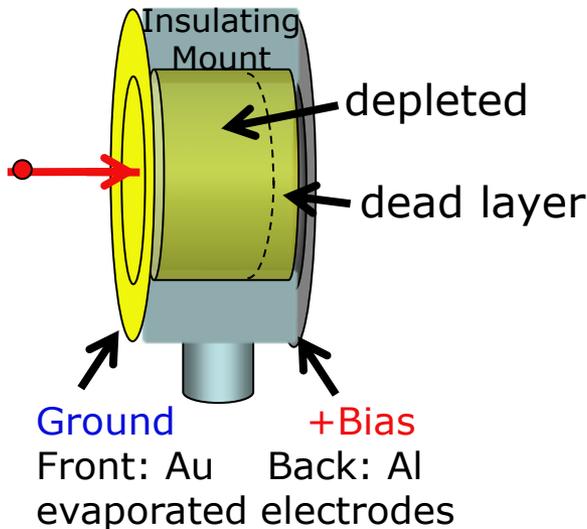
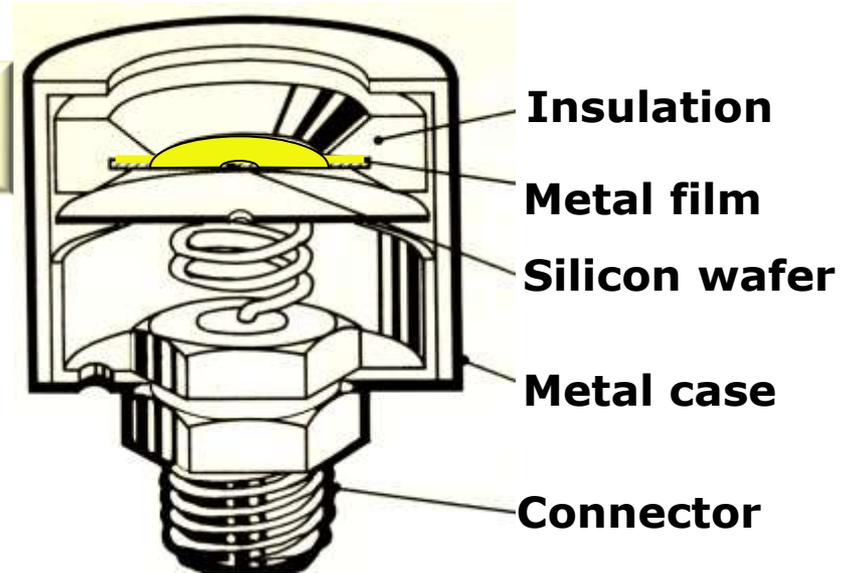
# Surface Barrier Detectors

Semiconductor/ Metal Junction



Different Fermi energies adjust to on contact. Thin metal film on Si surface produces space charge, an effective barrier (contact potential) and depleted zone free of carriers. Apply reverse bias to increase depletion depth.

ORTEC  
HI detector



Possible: electrical depletion depth  $\sim 100\mu$   
dead layer  $d_d \leq 1\mu$ ,  $V \sim 0.5V/\mu$ , "Over-bias" reduces  $d_d$

Used in medium-precision charged-particle spectroscopy ( $\alpha$  particles,  $\beta$  particles).

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Next:

# Gas Amplification Counters