

Agenda: TD of Dilute Gases, Simple Processes

- Fundamental ideal gas laws,
Equations of state (EoS)
Isothermal expansion/compression
Adiabatic expansion/compression
- Circular processes
Work, heat, and entropy in Carnot processes
Carnot, Kelvin/Clausius 2ndLTD
Thermodynamic driving potentials
Application chem reaction
- Equation of state of real gases
Van der Waals & other models
Activity, fugacity
Phase equilibria

Reading Assignments

Weeks 7 & 8

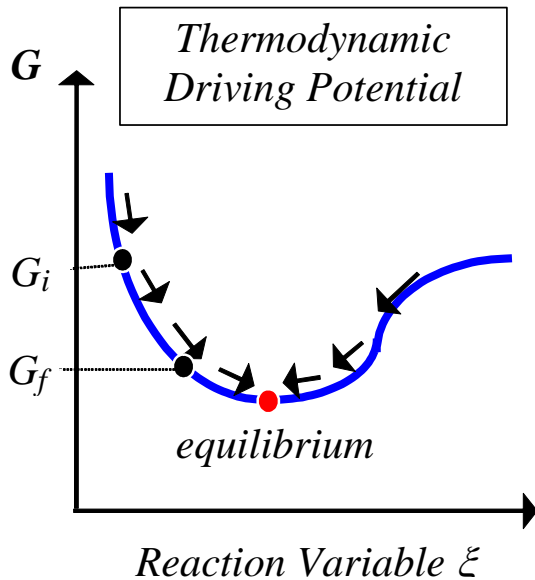
LN IV.4-7:

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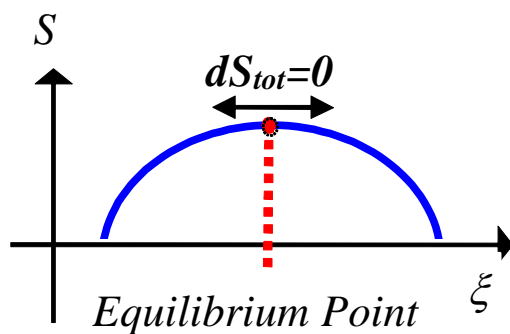
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- Statistical Mechanics
Quantum statistics
Partition functions

Reversible Work and Free Energy



Thermal Equilibrium

Stable Equilibrium



Maximum work a canonical system @ $S \leq S_{\max}$ can perform spontaneously using thermal energy dq

Net heat+work changes $\left\{ \begin{array}{l} \text{int. energy} \\ \text{enthalpy} \end{array} \right\} \rightarrow \left\{ (q+w) = \begin{array}{l} dU \\ dH \end{array} \right\}$

Helmholtz Free Energy

$$A = U - T \cdot S$$

$V, T = \text{const.}$

Work in reversible processes

$$dU = q_{\text{rev}} + w_{\text{rev}} = T \cdot dS + w_{\text{rev}}$$

$$w_{\text{rev}} = dU - T \cdot dS = d(U - T \cdot S) = dA$$

Gibbs Free Energy

$$G = H - T \cdot S$$

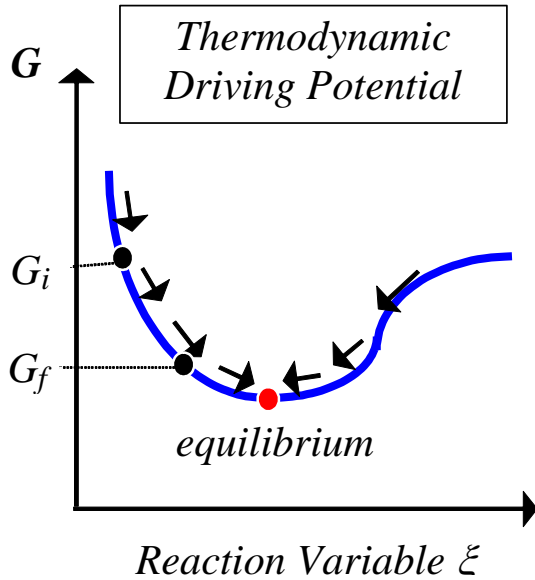
$P, T = \text{const.}$

Work in reversible processes

$$dH = q_{\text{rev}} + w_{\text{rev}} = T \cdot dS + w_{\text{rev}}$$

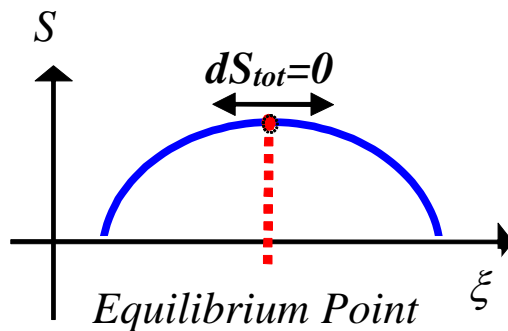
$$w_{\text{rev}} = dH - T \cdot dS = d(H - T \cdot S) = dG$$

Reversible Work and Free Energy



Thermal Equilibrium

Stable Equilibrium



Maximum work by a canonical system @ $S < S_{max}$

Gibbs Free Energy $G = H - T \cdot S$ $T = \text{const.}$

Work or G in reversible IG processes ?

$$w_{rev} = dH - T \cdot dS = d(H - T \cdot S) = dG$$

NOTE: $dH \neq q_{rev}$
since $p \neq \text{const}$

$$dG = d(U + p \cdot V) - d(T \cdot S)$$

At $T = \text{const}$

$$dG = dU + dp \cdot V + p \cdot dV - \cancel{dT} \cdot S - \underline{T \cdot dS}$$

with $\uparrow = q + w = \underline{T \cdot dS} - \underline{p \cdot dV}$

$$\rightarrow dG|_T = dp \cdot V \rightarrow \left(\frac{\partial G}{\partial p} \right)_T = V \leftarrow p = \text{dependence}$$



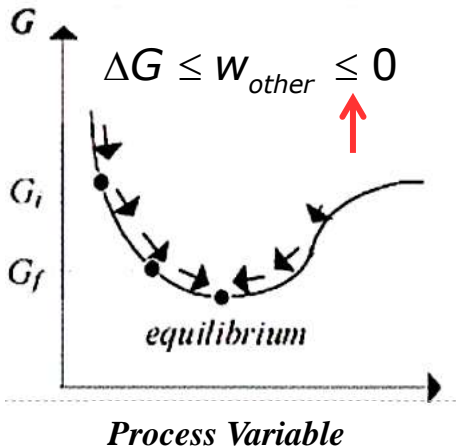
Work in isothermal reversible IG process

$$w_{rev} = dG|_T = V \cdot dp$$

Thermodynamic Energies & Driving Potentials

Potential energy gain in energetically allowed process (chem. rxn

$$\Delta H_{rxn}^0 = H_{products}^0 - H_{reagents}^0 < 0 \text{ does not guarantee spontaneity} \rightarrow \boxed{\Delta S > 0?}$$



Helmholtz free energy $A = U - T \cdot S$ used @ const V, T

System will do spontaneously : $w_{process} = \Delta A_{process} < 0$

Gibbs free energy

$G = H - T \cdot S$ used @ const p, T

System will do spontaneously : $w_{process} = \Delta G_{process} < 0$

Example : "extra" work a system can do beyond pV : $w_{other} = \Delta G$

Electrical potential Φ , charge $e \rightarrow$ $\boxed{\text{work} = e \cdot d\Phi}$

$$dG = d(U + e \cdot \Phi + p \cdot V) - d(T \cdot S) \quad (p, T = \text{const})$$

$$dG = \cancel{dU} + \underline{d(e \cdot \Phi)} + \cancel{dp \cdot V} + \cancel{p \cdot dV} - \cancel{dT \cdot S} - T \cdot dS$$

$$\uparrow dU = \text{heat} + \text{work on system} = \cancel{T \cdot dS} - \cancel{p \cdot dV}$$

Extra electrical work

$$dG = d(e \cdot \Phi) = w_{other}$$

Spontaneous if $dG < 0$

Canonical Thermodynamic IG Entropy Connections

Differentials of state functions **U, H, S, A, G, V, N** **$T=const$**

Ideal gas: $pV = Nk_B T$ and **$U = (3/2)Nk_B T$**

General 1st. LTD any process **$dU = dq - p \cdot dV + \mu \cdot dN + \dots$**

Reversible processes \rightarrow **$dU = T \cdot dS - p \cdot dV + \mu \cdot dN + \dots$**

$$dU(S, V, N) = \left(\frac{\partial U}{\partial S} \right)_{V, N, \dots} \cdot dS + \left(\frac{\partial U}{\partial V} \right)_{S, N, \dots} \cdot dV + \left(\frac{\partial U}{\partial N} \right)_{S, V, \dots} \cdot dN + \dots$$

$$\left(\frac{\partial U}{\partial S} \right)_{V, N} = T \quad \left(\frac{\partial U}{\partial V} \right)_{S, N} = -p \quad \left(\frac{\partial U}{\partial N} \right)_{S, V} = \mu \dots$$

Set $N=const$
 $\rightarrow dN=0$

Complete differential

$$dS = \frac{dU}{T} + \frac{p}{T} dV = \frac{3}{2} Nk_B \frac{dT}{T} + Nk_B \frac{dV}{V}$$

Integrate

$$S(T, V) = S_0(T_0, V_0) + Nk_B \cdot \ln \left\{ \left(\frac{T}{T_0} \right)^{3/2} \cdot \left(\frac{V}{V_0} \right) \right\} = S_0(T_0, V_0) + Nk_B \cdot \ln \left\{ \left(\frac{T}{T_0} \right)^{5/2} \cdot \left(\frac{p_0}{p} \right) \right\}$$

Since $V \propto 1/p$

Canonical Thermodynamic IG Entropy Connections

Differentials of state functions **U, H, S, A, G, V, N** **T=const.**

Ideal gas: $pV = Nk_B T$ and **U = (3/2)Nk_B T** per particle

General 1st. LTD any process **dU = dq - p·dV + μ·dN + ...**

Reversible processes → **dU = T·dS - p·dV + μ·dN + ...**

$$S(T, V) = S_0(T_0, V_0) + Nk_B \cdot \ln \left\{ \left(\frac{T}{T_0} \right)^{3/2} \cdot \left(\frac{V}{V_0} \right) \right\} = S_0(T_0, V_0) + Nk_B \cdot \ln \left\{ \left(\frac{T}{T_0} \right)^{5/2} \cdot \left(\frac{p_0}{p} \right) \right\}$$



Similarly

$$S(T, V, \dots) = \left(\frac{\partial S}{\partial T} \right)_V \cdot T + \left(\frac{\partial S}{\partial V} \right)_T \cdot V + \dots$$

$$\left(\frac{\partial S}{\partial T} \right)_V = \frac{1}{T} \left(\frac{\partial U}{\partial T} \right)_V = \frac{1}{T} C_V \quad \left(\frac{\partial S}{\partial V} \right)_T = \frac{1}{T} \left(\frac{\partial U}{\partial V} \right)_T + \frac{1}{T} p = \frac{1}{T} p$$

$$S(V, T) = C_V + \frac{1}{T} p \cdot V = C_V + k_B = \frac{3}{2} k_B + k_B = \frac{5}{2} k_B \quad \longrightarrow \quad S(V, T) = \frac{5}{2} k_B = C_P$$

Thermodynamic Relations: Free Energy

$$dA(V, T) = d(U - TS) = dU - SdT - TdS = (dq_{\text{rev}} - pdV) - SdT - TdS$$

$$\rightarrow dA(V, T) = -pdV - SdT$$

$$dA(V, T) = \left(\frac{\partial A}{\partial T}\right)_V dT + \left(\frac{\partial A}{\partial V}\right)_T dV \quad \left(\frac{\partial A}{\partial T}\right)_V = -S \quad \left(\frac{\partial A}{\partial V}\right)_T = -p$$

$$dG(p, T, N) = Vdp - SdT + \mu \cdot dN = \left(\frac{\partial G}{\partial p}\right)_{T, N} \cdot dp + \left(\frac{\partial G}{\partial T}\right)_{p, N} \cdot dT + \left(\frac{\partial G}{\partial N}\right)_{p, T} \cdot dN$$

$$\left(\frac{\partial G}{\partial p}\right)_{T, N} = V \quad \left(\frac{\partial G}{\partial T}\right)_{p, N} = -S \quad \left(\frac{\partial G}{\partial N}\right)_{p, T} = \mu$$

Gibbs-Duhem relation

$$d\mu = V \cdot dp - S \cdot dT$$

$$\uparrow k_B T / p$$

$$d\mu = k_B T \cdot d \ln(p) - \left(s_0 k_B + k_B \cdot \ln \left\{ \left(\frac{T}{T_0} \right)^{5/2} \right\} \right) dT$$

Thermodynamic Equilibrium State Functions

State Function	Fundamental Total Differential	@Equilibrium
$U(S, V, N)$	$dU = TdS - pdV + \sum_j \mu_j dN_j$ (mix, species j)	Minimum
$S(U, V, N)$	$dS = (1/T)dU + (p/T)dV - \sum_j (\mu_j/T)dN_j$	Maximum
$H(S, p, N) = U + pV$	$dH = TdS + Vdp + \sum_j \mu_j dN_j$	Minimum
$A(T, V, N) = U - TS$	$dA = -S dT - pdV + \sum_j \mu_j dN_j$	Minimum
$G(T, p, N) = H - TS$	$dG = -S dT + V dp + \sum_j \mu_j dN_j$	Minimum

Mixtures $N = \sum_j dN_j$

Chemical potential per particle type i : $\mu_i = \left(\frac{dG}{dN_i} \right)_{p,T}$

For canonical ensembles @

$V = \text{const.}, N_i = \text{const.}$ $dU = TdS - \cancel{pdV} + \cancel{\sum_j \mu_j dN_j} \rightarrow dU = TdS$

$P = \text{const.}, N_i = \text{const.}$ $dH = TdS + \cancel{Vdp} + \cancel{\sum_j \mu_j dN_j} \rightarrow dH = TdS$

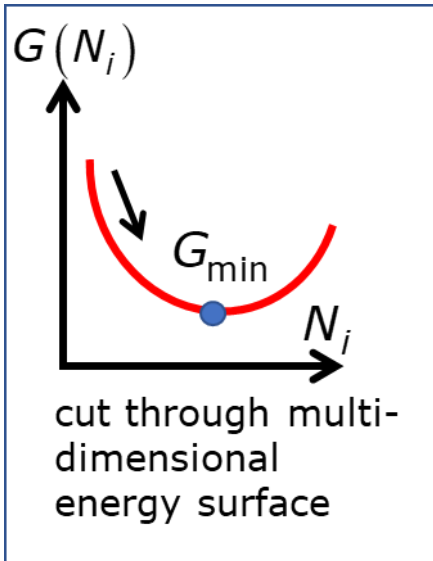
Reversible energy changes $dq = dU$ @ T, V and $dq = dH$ @ p, T

$$dS_{\text{rev}} = \frac{dq}{T}$$

Irreversible internal entropy changes (randomness)

$$dS_{\text{irrev}} > 0$$

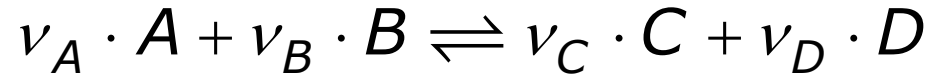
Grand Canonical PF: Application to Reactions



For EoS $p \cdot V = n \cdot R \cdot T$

$R = 8.31446262 \text{ J/K} \cdot \text{mol}$

Generic chemical reaction in gas phase



Change in chemical potential (@T) $\rightarrow \mu_i = \left(\frac{\partial G}{\partial N_i} \right)$

$$\Delta\mu = \nu_C \cdot \mu_C + \nu_D \cdot \mu_D - \nu_A \cdot \mu_A - \nu_B \cdot \mu_B \quad \nu_i \text{ in moles}$$

$$\mu_i(p_i, T) = \mu_i^0(T) + R \cdot T \cdot \ln(p_i/p_0); \quad p_0 := 101\text{kPa}$$

$$\Delta\mu = \Delta\mu^0 + R \cdot T \cdot \ln \left[\frac{(p_C)^{\nu_C} (p_D)^{\nu_D}}{(p_A)^{\nu_A} (p_B)^{\nu_B}} \right] \text{ all relative to } p_0$$

$$S = S_{max} \rightarrow \Delta\mu = 0 \rightarrow \Delta\mu^0 = -R \cdot T \cdot \ln \left[\frac{(p_C)^{\nu_C} (p_D)^{\nu_D}}{(p_A)^{\nu_A} (p_B)^{\nu_B}} \right] = -RT \cdot \ln K_p$$

Equilibrium Constant

$$K_p = \left[\frac{(p_C)^{\nu_C} (p_D)^{\nu_D}}{(p_A)^{\nu_A} (p_B)^{\nu_B}} \right] = \exp \left\{ -\frac{\Delta\mu^0}{R \cdot T} \right\}$$

Similar for $p \rightarrow (n/V) = [\]$

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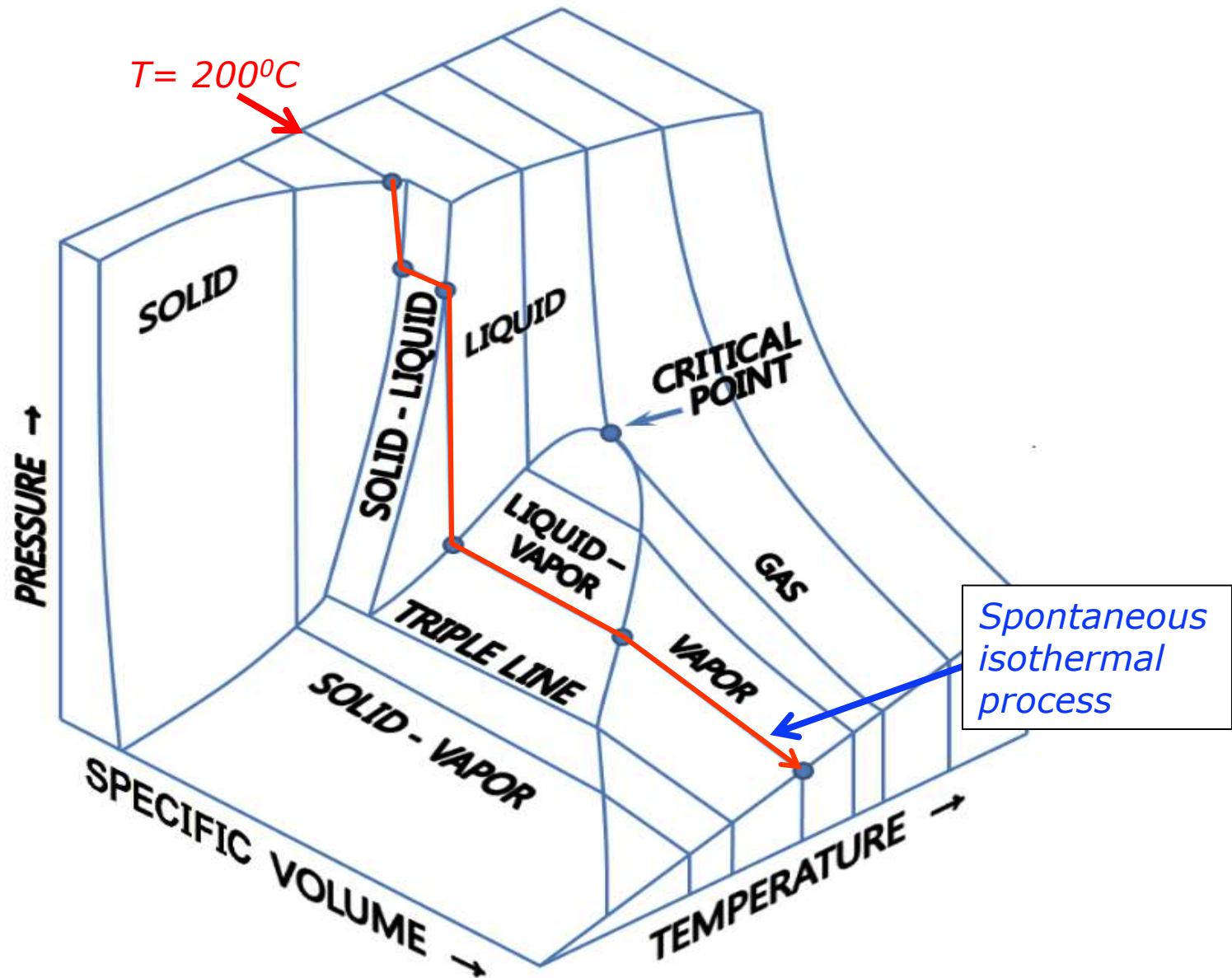
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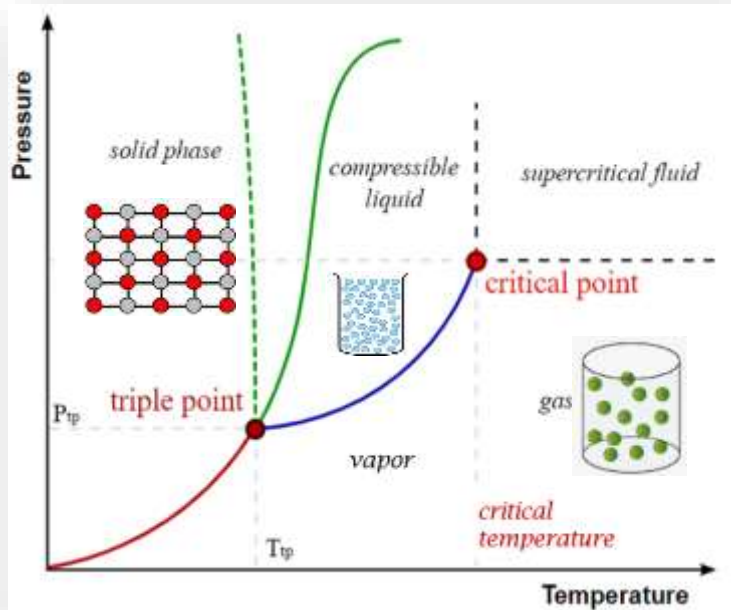
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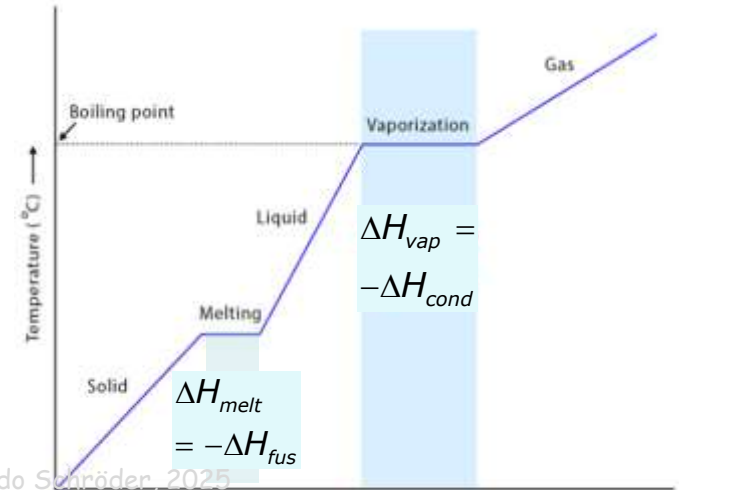
3D Phase Diagram for Water



Thermodynamic Energies in Phase Changes



Material	Formula	Critical pressure P_c		Critical temperature T_c		$k = C_p/C_v$
		psia	bar (abs)	°F	°C	
Water	H ₂ O	3206	221	705	374	1.32



Internal structural energy : 1mole substance
 extensive (additive) state function energy U
Enthalpy = Structural Energy +
+ $p - V$ work to access space

$$H = U + P \cdot V$$

→
 Ex. process @ $p = \text{const}$ → $\Delta H|_p = \Delta U + p \cdot dV$
 = extensive state function

Reference energy H^0 @ *standard state*
 $T = 25^\circ\text{C}$, $P = 1\text{bar}$ ← *large databases*
 $H^0 = 0$ for pure elemental substances

Phase changes $\Delta H \neq 0$
 → latent heat transfer $\Delta T = 0$

@ 1bar, $T = 100^\circ\text{C}$ (\neq the standard state)
 $\text{H}_2\text{O}(\ell) \rightarrow \text{H}_2\text{O}(\text{g}) \rightarrow \Delta H_{\text{vap}} = +40.7 \text{ kJ/mol}$
 $\text{H}_2\text{O}(\text{g}) \rightarrow \text{H}_2\text{O}(\ell) \rightarrow \Delta H_{\text{cond}} = -40.7 \text{ kJ/mol}$

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